

- 1 -

## Duistermaat- Heckman localisation formula and locus of vector fields.

12 October 2013 § 0

About two years ago (summer 2012) Dasha Belavin explained how to calculate an integral

$$Z(t) = \int e^{t d_K \omega} \quad (0.1)$$

( $\omega$ -1-form,  $d_K = d + (K)$ ). He showed first that this integral does not depend on  $t$ , then showed that it is localised at zeros of vector field  $K$ :

$$Z(H) \sim \frac{1}{\sqrt{\det \frac{\partial K}{\partial x}}} \Big|_{K=0} \quad (0.2)$$

It is typical localisation formula.

I tried to revive these calculations. On one hand they are leading to Duistermaat- Heckman formula in more less general case.  
On the other hand we may discuss what is interesting to analyze geometrical meaning of answer.

### § 1.

localisation

Two words about Duistermaat- Heckman formula (DHL)-formula.

Let  $M$  be compact manifold

$(M^{2n}, \Omega)$  be compact symplectic manifold

- 2 -

Let  $H$  be an Hamiltonian, such that  
the vector field

$$K = D_H : \Omega \rightarrow D_H = -dH$$

is a compact vector field

(i.e. it generates compact subgroup  $e^{tk}$   
in the group of diffeomorphisms)

Then

$\int \Omega^n e^{iH}$  is localized at zero locus  
~~of vector field K~~  
of vector field  $K$

and  $\int \Omega^n e^{iH} = \sum_i e^{\frac{iH(x_i)}{\sqrt{\det \text{Hess } H(x_i)}}}$

$$\int \Omega^n e^{iH} = \sum_{x_i: K(x_i)=0} e^{\frac{iH(x_i)}{\sqrt{\det \text{Hess } H(x_i)}}}$$

(we suppose that  $K(x_i)$  are not-degenerate).

This is famous Duistermaat-Heckman formula.

We will consider here a  
special but very illuminating case  
of this formula.

etude

[see in more detail the next file].

- 3 -

We consider now the following set up:

Let  $w$  be 1-form on  $M$  ( $\dim M = 2m$ )

such that  $\Omega = dw$  defines symplectic structure.

(of course condition  $\Omega = dw$  is in contradiction with

connectedness of  $M$ :  $\int \Omega^n \neq 0$ , but we ignore now this.)

E.g. we suppose that  $M$  is not compact

Let  $K$  be a vector field such that

$$L_K w = d(w \lrcorner K) + d(w \lrcorner K) = 0$$

Then it is evident that  $K$  is

Hamiltonian vector field of  $H = w \lrcorner K$

$$\Omega \lrcorner K = dw \lrcorner K = -d(w \lrcorner K) = -dH.$$

$$\begin{array}{c} w \\ \cancel{\Omega = d\Omega} / \quad \cancel{H = w \lrcorner K} \quad d_K w = \\ H = -dH = (d + L_K)w = \\ H_+ = \Omega + H. \end{array}$$

We see that

$$\int_{\mathbb{D}^n} e^{iH} = \int e^{iH + i\Omega}$$

$$= \int e^{id_K w}$$

We come to integral (o)

§ 2

- 4 -

Calculation of  $\int e^{it \Delta_K w}.$

Consider

$$Z(t) = \int_M e^{it \Delta_K w}. \quad [\Delta_K^2 = \Delta_K]$$

Show that  $Z(t)$  does not depend on  $t$ .

$$\begin{aligned} \frac{dZ(t)}{dt} &= i \int_M \Delta_K w e^{it \Delta_K w} = \\ &= i \int_M \Delta_K (w e^{it \Delta_K w}) = i \int_M (\Delta_K w e^{it \Delta_K w}) = 0 \end{aligned} \quad (2.1)$$

(under some technical conditions).

[  $\int_L \Delta_K w = 0$  since form  $L_K w$  has rank  $\leq 2n-1$  ]

We see that  $Z(t)$  does not depend on  $t$ .  
Hence we can calculate  $Z(t)$  at  $t \rightarrow \infty$ .

$$\begin{aligned} \int e^{it \Delta_K w} &= \int e^{it(\Delta + H)} = \\ &= d\omega = \Omega, \quad w \perp K = H. \quad (\Delta_K w = 0) \\ &= \sum \frac{i^{mt^n}}{n!} \int \Omega^n e^{itH} = \frac{i^{mt^m}}{m!} \int \Omega^m e^{itH} \end{aligned} \quad (2.2)$$

(dim  $M = 2m$ )

Calculate using stationary phase method:

$$dH = d(w \perp K) = -d\omega \perp K \quad (2.3)$$

Locus of  $dH =$  locus of  $K$

§ 2 - 5 -

We see that at stationary point  $dH = 0$   
 Hermitian is:

$$\frac{\partial^2 H}{\partial x^i \partial x^K} \Big|_{K=0} = \frac{\partial}{\partial x^i} \frac{\partial}{\partial x^K} (W_K K^i) \Big|_{K=0} = \frac{\partial}{\partial x^i} (\Omega_{ip} K^p) =$$

$$(d\langle e | S | k \rangle = - d\langle e | S | k \rangle)$$

$$= \underline{\Omega_{ip}} \frac{\partial K^p}{\partial x^i} \quad (\underline{H(x_0) \geq 0 \text{ for } K(x_0) \neq 0})$$

Hence:

$$\int \Omega^n e^{i\theta H} \sim \underbrace{\sqrt{\det \Omega} (e^{i\theta H(x_0)})}_{\sqrt{\det(\Omega \cdot \frac{\partial K}{\partial x})}}$$

$$\sim \sum_{x_0} \frac{e^{i\theta H(x_0)}}{\sqrt{\det \frac{\partial K}{\partial x}}}$$

Note:  $\frac{\partial K}{\partial x}$  is linear operator at particular  $x_0$

$$L_K = \frac{\partial K}{\partial x}, \quad [L_K u] = -[K, u].$$

We see that answer does not depend on  
 choice of  $\underline{W}$ .

$$\int \sim \frac{1}{\sqrt{\det \frac{\partial K}{\partial x}}}$$

Our formula is a special case of DHL formula (In particular  $H(x_i) = 0$ ).

On the other hand this formula emphasizes the role of vector field  $K$ : It states that

$$\int e^{it(dW + L_K(x))} \underset{x \in K(0)}{\int} \frac{dx}{\sqrt{\det \frac{\partial K}{\partial x}}}$$

depends only on  $K$  at least in the case if  $W$  is an "arbitrary"  $K$ -invariant 1-form (of course  $dW$  is not-degenerate).

It is useful to study DHL formula in its supersymmetric manifestation.

1. A. Nersessian, "An fibrackets and realisation of path integral,"

(See 2 A. Schwarz, O. Zaboronsky

"Supersymmetry and localization"

JETP Lett. 58.1 (1993)

CMP (1995 or 96)

(See for detail  
next etude)

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