

# Higher order Koszul brackets

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*The talk is based on the work with Ted Voronov*

# Contents

Abstracts

Poisson manifold and....

Higher brackets

## Papers that talk is based on are

- [1] H.M.Khudaverdian, Th. Voronov *Higher Poisson brackets and differential forms*, 2008a In: Geometric Methods in Physics. AIP Conference Proceedings 1079, American Institute of Physics, Melville, New York, 2008, 203-215., arXiv: 0808.3406
- [2] Th. Voronov, *Nonlinear pullback on functions and a formal category extending the category of supermanifolds*], arXiv: 1409.6475
- [3] Th. Voronov, *Microformal geometry*, arXiv: 1411.6720

## Abstract...

For an arbitrary manifold  $M$ , we consider supermanifolds  $\Pi TM$  and  $\Pi T^*M$ , where  $\Pi$  is the parity reversion functor. The space  $\Pi T^*M$  possesses canonical odd Schouten bracket and space  $\Pi TM$  possesses canonical de Rham differential  $d$ . An arbitrary even function  $P$  on  $\Pi T^*M$  such that  $[P, P] = 0$  induces a homotopy Poisson bracket on  $M$ , a differential,  $d_P$  on  $\Pi T^*M$ , and higher Koszul brackets on  $\Pi TM$ . (If  $P$  is fiberwise quadratic, then we arrive at standard structures of Poisson geometry.) Using the language of  $Q$ -manifolds and in particular of Lie algebroids, we study the interplay between canonical structures and structures depending on  $P$ . Then using just recently invented theory of thick morphisms we construct a non-linear map between the  $L_\infty$  algebra of functions on  $\Pi TM$  with higher Koszul brackets and the Lie algebra of functions on  $\Pi T^*M$  with the canonical odd Schouten bracket.



## Poisson manifold

Let  $M$  be Poisson manifold with Poisson tensor  $P = P^{ab} \partial_b \wedge \partial_a$

$$\{f, g\} = \{f, g\}_P = \frac{\partial f}{\partial x^a} P^{ab} \frac{\partial g}{\partial x^b}.$$

$$\{\{f, g\}, h\} + \{\{g, h\}, f\} + \{\{h, f\}, g\} = 0,$$



$$P^{ar} \partial_r P^{bc} + P^{br} \partial_r P^{ca} + P^{cr} \partial_r P^{ab} = 0.$$

If  $P$  is non-degenerate, then  $\omega = (P^{-1})_{ab} dx^a \wedge dx^b$  is closed non-degenerate form defining symplectic structure on  $M$ .

## Differentials

$d$ —de Rham differential,  $d: \Omega^k(M) \rightarrow \Omega^{k+1}(M)$ ,

$$d^2 = 0, df = \frac{\partial f}{\partial x^a} dx^a, \quad d(\omega \wedge \rho) = d\omega \wedge \rho + (-1)^{\rho(\omega)} \omega \wedge d\rho$$

$d_P$ —Lichnerowicz- Poisson differential,  $d_P: \mathfrak{X}^k(M) \rightarrow \mathfrak{X}^{k+1}(M)$ ,

$$d_P^2 = 0, df = \frac{\partial f}{\partial x^b} P^{ba} \frac{\partial}{\partial x^a}$$

$d_P P = 0 \leftrightarrow$  Jacobi identity for Poisson bracket  $\{, \}$

## Differential forms and multivector fields

$\mathfrak{A}^*$  space multivector fields on  $M$ ,

$\Omega^*$  space of differential forms on  $M$ ,

$$\begin{array}{ccc} \mathfrak{A}^k(M) & \xrightarrow{d_P} & \mathfrak{A}^{k+1}(M) \\ \uparrow & & \uparrow \\ \Omega^k(M) & \xrightarrow{d} & \Omega^{k+1}(M) \end{array}$$



## Differential forms and multivector fields

$\mathfrak{A}^*$  — multivector fields on  $M$  = functions on  $\Pi T^*M$

$\Omega^*$  — differential forms on  $M$  = functions on  $\Pi TM$ ,

$$\begin{array}{ccc}
 \mathfrak{A}^k(M) & \xrightarrow{d_P} & \mathfrak{A}^{k+1}(M) & & C(\Pi T^*M) & \xrightarrow{d_P} & C(\Pi T^*M) \\
 \uparrow & & \uparrow & & \uparrow & & \uparrow \\
 \Omega^k(M) & \xrightarrow{d} & \Omega^{k+1}(M) & & C(\Pi TM) & \xrightarrow{d} & C(\Pi TM)
 \end{array}$$

$$d\omega(x, \xi) = \xi^a \frac{\partial}{\partial x^a} \omega(x, \xi), \quad d_P F(x, \theta) = (P, F)_1,$$

$(P, F)_1$ -canonical odd Poisson bracket on  $\Pi T^*M$ .

$x^a = (x^1, \dots, x^n)$  — coordinates on  $M$

$(x^a, \xi^b) = (x^1, \dots, x^n; \xi^1, \dots, \xi^n)$ , —coordinates on  $\Pi TM$

$$\rho(\xi^a) = \rho(x^a) + 1, x^{a'} = x^{a'}(x^a) \rightarrow \xi^{a'} = \xi^a \frac{\partial x^{a'}}{\partial x^a}. \quad (dx^a \leftrightarrow \xi^a).$$

‘ Respectively

$(x^a, \theta_b) = (x^1, \dots, x^n; \theta_1, \dots, \theta_n)$ , —coordinates on  $\Pi T^*M$

$$\rho(\theta_a) = \rho(x^a) + 1, x^{a'} = x^{a'}(x^a) \rightarrow \theta_{a'} = \theta_a \frac{\partial x^a}{\partial x^{a'}}. \quad (\partial_a \leftrightarrow \theta_a).$$

‘

## Example

$$\Omega^* \ni \omega = l_a dx^a + r_{ab} dx^a \wedge dx^b \leftrightarrow \omega(x, \xi) = l_a \xi^a + r_{ab} \xi^a \xi^b \in C(\Pi TM)$$

$$\mathfrak{A}^* \ni F = X^a \partial_a + M^{ab} \partial_a \wedge \partial_b \leftrightarrow F(x, \theta) = X^a \theta_a + M^{ab} \theta_a \theta_b \in C(\Pi T^*M).$$

## Canonical odd Poisson bracket

$F, G$  multivector fields

$[F, G]$  Schouten commutator'

$$\mathbf{X} = X^a \partial_a, [\mathbf{X}, F] = \mathfrak{L}_{\mathbf{X}} F$$

$$P = P^{ab} \partial_a \wedge \partial_b, [P, F] = d_P F,$$

$F, G$  functions on  $\Pi T^*M$

$[F, G]$  odd Poisson bracket'

$$[\mathbf{X}, F] = (X^a \theta_a, F(x, \theta))$$

$$d_P F = (P, F) = (P^{ab} \theta_a \theta_b, F(x, \theta))_1$$

$$[F(x, \theta), G(x, \theta)] = \frac{\partial F(x, \theta)}{\partial x^a} \frac{\partial G(x, \theta)}{\partial \theta_a} + (-1)^{p(F)} \frac{\partial F(x, \theta)}{\partial \theta_a} \frac{\partial G(x, \theta)}{\partial x^a}.$$

odd Poisson bracket

Schouten bracket

Buttin bracket

anti-bracket

Names are

## Koszul bracket on differential forms

$$\varphi_P^*: \begin{array}{c} C(\Pi T^*M) \\ \uparrow \\ \mathbf{C}(\Pi TM) \end{array} \quad \xi^a = P^{ab} \theta_b \text{ or } dx^a = P^b \partial_b$$

From  $\{, \}$  on functions to Koszul bracket on differential forms

$$[\omega, \sigma]_P = (\varphi_P^*)^{-1} ([\varphi_P^*(\omega), \varphi_P^*(\sigma)]_P).$$

$$[f, g]_P = 0, [f, dg]_P = (-1)^{p(f)} \{f, g\}_P, [df, dg]_P = (-1)^{p(f)} d(\{f, g\}_P)$$

This formula survives the limit if  $P$  is degenerate.

## Lie algebroid

$E \rightarrow M$ —vector bundle,

$[[, ]]$ —commutator on sections,  $\rho: E \rightarrow TM$ —anchor

$$[[\mathbf{s}_1(x), f(x)\mathbf{s}_2(x)]] = f(x)[[\mathbf{s}_1(x), \mathbf{s}_2(x)]] + \left(\rho(\widehat{\mathbf{s}_1(x)})f(x)\right)\mathbf{s}_2(x),$$

Jacobi identity:

$$[[[\mathbf{s}_1, \mathbf{s}_2]], \mathbf{s}_3]] + \text{cyclic permutations} = 0.$$

$$\mathbf{s}(x) = s^i(x)\mathbf{e}_i(x), \quad [[\mathbf{e}_i(x), \mathbf{e}_k(x)]] = c_{ik}^m(x)\mathbf{e}_m(x), \quad \rho(\mathbf{e}_i) = \rho_i^\mu \partial_\mu,$$

$$[[\mathbf{s}_1(x), \mathbf{s}_2(x)]] = \left( s_1^i s_2^k c_{ik}^m + s_1^i \rho_i^\mu \partial_\mu s_2^m(x) - s_2^i \rho_i^\mu \partial_\mu s_1^m(x) \right) \mathbf{e}_m$$

## Trivial examples of Lie algebroid

$\mathcal{G}$  — Lie algebra,  $\begin{array}{c} \mathcal{G} \\ \downarrow \\ * \end{array}$ , where  $[[, ]]$  — usual commutator,

tangent bundle  $\begin{array}{c} TM \\ \downarrow \\ M \end{array}$ , where  $[[, ]]$  — commutator of vector fields

For  $TM$  anchor is identity map

## Poisson algebroid

$(M, P)$  Poisson manifold,  $(P = P^{ab} \partial_b \wedge \partial_a, \{f, g\} = \partial_a f P^{ab} \partial_b g)$

$T^*M$

$\downarrow$ ,  $[[df, dg]] = d\{f, g\}$ , anchor  $\rho: \rho(\omega_a dx^a) = D_\omega = P^{ab} \omega_b \frac{\partial}{\partial x^a}$ ,  
 $M$

$$[[\omega_a dx^a, \sigma_b dx^b]] = \left( \frac{1}{2} \omega_a \sigma_b \partial_c P^{ab} + P^{ab} \omega_b \partial_a \sigma_c - (\omega \leftrightarrow \sigma) \right) dx^c$$

(This is Koszul bracket  $[\cdot, \cdot]_P$  on 1-forms).

## Anchor—morphism of algebroids

$$\text{Anchor } \rho : \begin{pmatrix} T^*M \\ \downarrow \\ M \end{pmatrix} \rightarrow \begin{pmatrix} TM \\ \downarrow \\ M \end{pmatrix},$$

morphism of algebroid  $T^*M$  to tangent algebroid.

$$\rho[[\omega, \sigma]] = [\rho(\omega), \rho(\sigma)].$$



## One very useful object— $Q$ manifold

### Definition

A pair  $(M, Q)$  where  $M$  is (super)manifold, and  $Q$  is odd vector field on it such that

$$Q^2 = \frac{1}{2}[Q, Q] = 0$$

is called  $Q$ -manifold.

$Q$  is called homological vector field.

## Lie algebroid and its neighbours

Algebroid has different manifestations

$$\begin{array}{ccc}
 \Pi E & & E \\
 \downarrow & & \downarrow \\
 M & , & M
 \end{array}$$

$\Pi E$  is  $Q$  manifold with  $Q = \xi^k \xi^i c_{ik}^m \frac{\partial}{\partial \xi^m} + \xi^i \rho_i^\mu \frac{\partial}{\partial x^\mu}$

$E \rightarrow M$  is Lie algebroid with  $[[\mathbf{e}_i, \mathbf{e}_k]] = c_{ik}^m, \rho(\mathbf{e}_i) = \rho_i^\mu \frac{\partial}{\partial x^\mu}$

$$\begin{array}{ccc}
 E^* & \Pi E^* & \\
 \downarrow , & \downarrow & \text{---(even, odd)Poisson manifolds} \\
 M & M &
 \end{array}$$

Lie–Poisson bracket:

$$\{u_i, u_k\} = c_{ik}^m u_m, \{x^\mu, u_i\} = \rho_i^\mu, \{x^\mu, x^\nu\} = 0.$$

Neighbours of  $\mathcal{G} \rightarrow *$ 

$$\begin{array}{ccc}
 \begin{array}{c} \Pi\mathcal{G} \\ \downarrow \\ * \end{array} & & \begin{array}{c} \mathcal{G} \\ \downarrow \\ * \end{array} & & \begin{array}{c} \mathcal{G}^* \\ \downarrow \\ * \end{array} \\
 \underbrace{Q = \xi^i \xi^k c_{ik}^m \frac{\partial}{\partial \xi^m}}_{\text{homological vector field}} & , & \underbrace{[\mathbf{e}_i, \mathbf{e}_k] = c_{ik}^m \mathbf{e}_m}_{\text{structure constants}} & , & \underbrace{\{u_i, u_k\} = c_{ik}^m u_m}_{\text{Lie-Poisson bracket}}
 \end{array}$$

Neighbours of tangent algebroid  $TM \rightarrow M$ 

$$\begin{array}{c} \Pi TM \\ \downarrow \\ M \\ \underbrace{Q = \xi^m \frac{\partial}{\partial x^m}} \end{array},$$

homological vector field—de Rham differential  $d$   
(functions on  $\Pi TM$ )—differential forms on  $M$ )

 $T^*M$ 
 $\downarrow$ 
 $M$ 

canonical symplectic structure

 $\Pi T^*M$ 
 $\downarrow$ 
 $M$ 

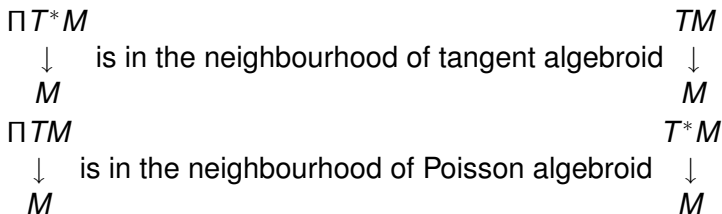
canonical odd symplectic structure

Neighbours of Poisson algebroid  $T^*M \rightarrow M$  $(M, P)$ —Poisson manifold,  $\{x^a, x^b\} = P^{ab}$ 

$$\begin{array}{ccc}
 \Pi T^*M & & T^*M \\
 \downarrow & & \downarrow \\
 M & & M \\
 \underbrace{Q = \theta_a \theta_b \frac{\partial P^{ba}}{\partial x^c} \frac{\partial}{\partial \theta_c} + \theta_a P^{ab} \frac{\partial}{\partial x^b}}_{\text{homological vector field}}, & & \text{Poisson algebroid} \\
 & & [[dx^a, dx^b]] = dP^{ab}, \rho(dx^a) = P^{ab} \partial_b
 \end{array}$$

$$\begin{array}{c}
 \Pi TM \\
 \downarrow \\
 M
 \end{array}$$

 $\{, \} = [, ]_P$  is Koszul bracket on  $\Pi TM$ .



$\underbrace{\Pi T^*M}$                        $\rightarrow$                        $\underbrace{\Pi TM}$   
 Odd canonical Poisson bracket      Odd Koszul bracket

i

$$\text{Linear map } \xi^a = \frac{1}{2} \frac{\partial P(x, \theta)}{\partial \theta_a} = P^{ab} \theta_b, \quad (dx^a = P^{ab} \partial_b) \quad )$$

## Question

What happens if even function  $P = P^{ab}(x, \theta)\theta_a\theta_b$  is replaced by an arbitrary even function  $P = P(x, \theta)$  which obeys the master-equation

$$[P, P] = 2 \frac{P(x, \theta)}{\partial x^a} \frac{P(x, \theta)}{\partial \theta^a} = 0.$$

(In the case  $P = P^{ab}(x, \theta)\theta_a\theta_b$  master-equation is just Jacobi identity for Poisson bracket  $\{, \}_P$  on  $M$ .)

## Higher Poisson brackets on $M$

$P$ :  $[P, P] = 0$  define higher brackets

$$\{f_1, f_2, \dots, f_n\}_P = [\dots [P, f_1], \dots, f_p] \Big|_M, \quad \Big|_M = \Big|_{\theta=0}.$$

$$P = P^a \theta_a + P^{ab} \theta_b \theta_a + P^{abc} \theta_c \theta_b \theta_a + \dots$$

$$\{X^a\}_P = P^a, \quad \{X^a, X^b\} = P^{ab}, \quad \{X^a, X^b, X^c\} = P^{abc} \dots$$



## From $\Pi T^*M$ to $\Pi TM$

$$C(\Pi T^*M) \rightarrow \mathfrak{X}(\Pi T^*M) \rightarrow C(T^*(\Pi T^*M)) \rightarrow C(T^*(\Pi T^*M))$$

Function  $P(x, \theta) \rightarrow$  Hamiltonian vector field  $D_F \rightarrow$   
 $\rightarrow$  Hamiltonian in  $T^*(\Pi T^*M) \rightarrow T^*(\Pi T^*M)$

The last map is Mackenzie Xu symplectomorphism

$$C(\Pi TM) \ni P = P(x, \theta) \rightarrow K = K_P(x, \xi) \in T^*(\Pi T^*M)$$

$$K_P(x, \xi, \rho, \pi) = \left( \rho_a \frac{\partial}{\partial \theta_a} P(x, \theta) + \xi^a \frac{\partial}{\partial x^a} P(x, \theta) \right) \Big|_{\theta \rightarrow \pi}$$

$(x^a, \xi^b | \rho_a, \pi_b)$  coordinates on  $T^*(\Pi TM)$ .

## Higher Koszul brackets on $M$

$P \in \Pi T^*M$  induces homotopy Poisson bracket in  $M$ ,  
 $K_P \in T^*(\Pi TM)$  induces homotopy odd Poisson bracket (higher Koszul bracket) on  $\Pi M$ ,

$$\{F_1, F_2, \dots, F_n\}_{K_P} = [\dots [K_P, F_1], \dots, F_n] \Big|_{\Pi M}, \quad \Big|_{\Pi M} = \Big|_{\rho=\pi=0}.$$

$$F = F(x, \xi) = f(x) + \xi^a f_a(x) + \dots, \quad (df = \xi^a \partial_a f),$$

$$[f]_P = 0, [f_1, f_2, \dots, f_k]_P = 0$$

$$[f_1, df_2, \dots, df_n] = f_1, f_2, \dots, f_n,$$

$$[df_1, df_2, \dots, df_n] = df_1, f_2, \dots, f_n,$$

## Q-manifolds

$$C(\pi T^*M) \xleftarrow{\text{morphism of Q-manif.}} C(\pi TM)$$

 $\pi T^*M$  $\pi TM$ 

Lichnerowicz Poisson differential  $d_P \rightarrow$  de Rham differential  
 Odd Poisson canonical bracket      Odd Koszul bracket

$$d = \xi^a \partial_a,$$

$$d_P: d_P f = [P, F], \quad d_P = \frac{\partial P}{\partial x^a} \frac{\partial}{\partial \theta_a} + \frac{\partial P}{\partial \theta_a} \frac{\partial}{\partial x^a}$$

If  $P = P^{ab}$  then the map

$$\Pi T^*M \rightarrow \Pi TM: \quad \xi^a = \frac{\partial P}{\partial \theta^a} = P^{ab}(x)\theta_b,$$

is linear in fibres. Morphism of  $Q$ -manifolds

$$C(\Pi T^*M) \leftarrow C(\Pi TM)$$

is its pull-back.

These linear maps interwin differentials  $d$  and  $d_P$ , their Hamiltonians, and their homological vector fields on infinite-dimensional spaces of functions.

It is more tricky if  $P(x, \theta)$  is an arbitrary function (solution of master-equation  $[S, S] = 0$ ). The map

$$\Pi T^*M \rightarrow \Pi TM: \quad \xi^a = \frac{\partial P}{\partial \theta^a} = P^{ab}(x)\theta_b,$$

and its pull-back is in general non-linear map.

$$\begin{array}{ccc} \Pi T^*M & \xrightarrow{\text{non-linear}} & \Pi TM \\ \Pi TM & \xleftarrow{\text{thick}} & \Pi T^*M \end{array}$$

i.e.

$$C(\Pi TM) \text{ non-linear map } C(\Pi T^*M)$$

This non-linear map defines morphism of  $Q$ -manifolds.

## Papers that talk is based on

- [1] H.M.Khudaverdian, Th. Voronov *Higher Poisson brackets and differential forms*, 2008a In: Geometric Methods in Physics. AIP Conference Proceedings 1079, American Institute of Physics, Melville, New York, 2008, 203-215., arXiv: 0808.3406
- [2] Th. Voronov, *Nonlinear pullback on functions and a formal category extending the category of supermanifolds*], arXiv: 1409.6475
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