

27 September 2018
First lecture.

(1)

States and Observables

Unitary space.

Complex vector space \mathcal{H} equipped with
scalar product $\langle \cdot, \cdot \rangle$
(hermitian positive definite form)

$$\langle \bar{x}, \bar{y} \rangle = \overline{\langle y, x \rangle}$$

$$(1) \langle \lambda x, y \rangle = \lambda \langle x, y \rangle \quad \left(\text{In physics sometimes is: } \langle x, \lambda y \rangle = \lambda \langle x, y \rangle \right)$$
$$(2) \langle x, \lambda y \rangle = \bar{\lambda} \langle x, y \rangle$$

Exercise. Deduce (1) \rightarrow (2)

Remark (in physics use $\langle \lambda x, y \rangle = \bar{\lambda} \langle x, y \rangle$)

Example

$$\mathbb{C}^2 = \left\{ \begin{pmatrix} a \\ b \end{pmatrix}, a, b \in \mathbb{C} \right\}$$

$$\left\langle \begin{pmatrix} a \\ b \end{pmatrix}, \begin{pmatrix} a' \\ b' \end{pmatrix} \right\rangle = a \bar{a}' + b \bar{b}'$$

$$\begin{pmatrix} a \\ b \end{pmatrix} = a \uparrow + b \downarrow \quad \left(\uparrow = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \downarrow = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \right)$$

Sometimes \mathbb{C}^2 is called the space of spinors

Example 2

$$V = \left\{ f: f \in C^\infty(\mathbb{R}^3), \text{ such that } \int f \bar{f} dx < \infty \right\}$$

$$\langle f, g \rangle = \int f(x) \bar{g(x)} d^3x$$

Exercise Show that this is scalar product

Remark The space $(C^\infty(\mathbb{R}^3), \langle \cdot, \cdot \rangle)$ is not complete. Mathematicians say that this is so called pre-Hilbert space. (2)

Hilbert space is a completion of pre-Hilbert space ~~with respect to~~ i.e. ~~see~~ with respect to metric induced by scalar product.

Recall it:

We say that $\{x_n\}$ is Cauchy sequence if $\lim_{n, m \rightarrow \infty} \|x_n - x_m\| = 0$ ($\|\vec{x}\| = \sqrt{\langle \vec{x}, \vec{x} \rangle}$)

pre-Hilbert space \mathcal{H} is Hilbert if it is closed, i.e. every Cauchy sequence has limit

$$\overline{C^\infty(\mathbb{R}^n), \langle \cdot, \cdot \rangle} = L^2(\mathbb{R}^n)$$

Physicists usually do not care.

Exercise. Let f be a real function. Show that $\langle f, f \rangle < \infty$.

Exercise

$$\forall f, \int f^2 dx < \infty \Rightarrow \forall f, g \int fg < \infty$$

or in other words

Does scalar product exist

To answer this question consider

Cauchy - Bunyakovsky - Schwarz inequality

< , >

$$\| \langle f, g \rangle \|^2 \leq \| f \|^2 \| g \|^2$$

Proof.

It is instructive first to prove for real case.

$$P(x) = \langle x f + g, x f + g \rangle \geq 0$$

↓
CBS.

Complex case

$$P(z) = \langle z f + g, z f + g \rangle =$$

$$= (x^2 + y^2) A + x b - y \beta + C. =$$

$$A = \langle f, f \rangle \quad b + i\beta = \langle f, g \rangle, \quad C = \langle g, g \rangle$$

$z = x + iy$

$$= A \left(x + \frac{b}{A} \right)^2 + A \left(y - \frac{\beta}{A} \right)^2 + C - \frac{b^2}{A} - \frac{\beta^2}{A}.$$

(the one $A=0 \Rightarrow f=0$ evidently)

$$CA > b^2 + \beta^2 \quad \square$$

(Details on tutorial)

$$f, g \in \mathcal{H} \Leftrightarrow \int f \bar{f} d\alpha < \infty, \int g \bar{g} < \infty$$

↓ CBS

$$\sqrt{\langle f, g \rangle} < \langle f, 1 \rangle \langle 1, g \rangle < \infty$$

Remark - \mathcal{H} - is called sometimes
Hilbert space.

(4)

R. We avoid discuss. of difference
between finite and inf. dim. case.

Space of states in QM

- $(\mathcal{H}, \langle, \rangle)$ - unitary space
(Hilbert space)

State (pure state) -

one-dimensional complex subspace

$$[\psi], \psi \in \mathcal{H}, \psi \neq 0.$$

ψ is called ψ -function.

Superposition of states.

$$\psi = C_1 \psi_1 + \dots + C_n \psi_n.$$

Postulat: ψ -space of states is complex
linear space (superposition postulate).

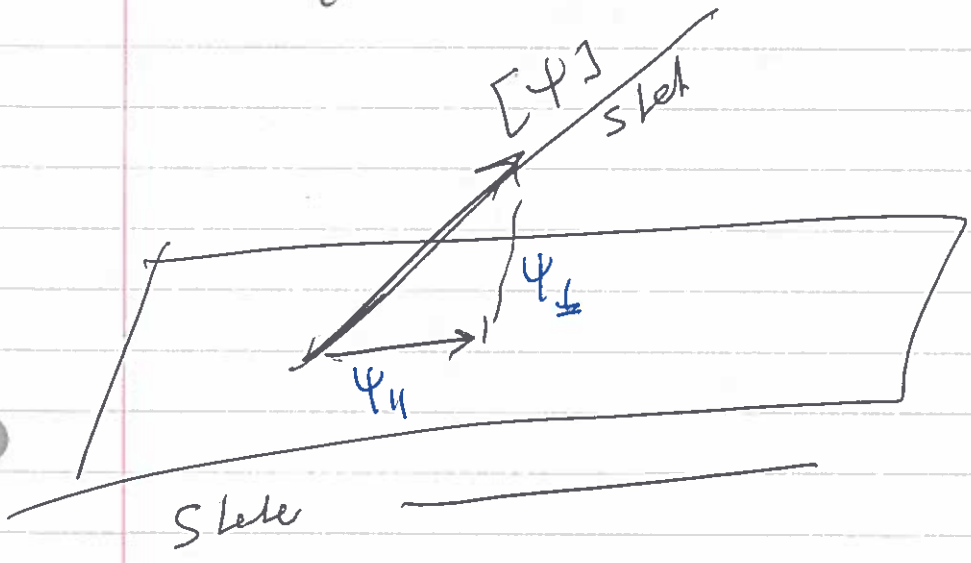
Measurement

In classical mechanics
 state $\xrightarrow{\text{Measurement (Experiment)}}$ "point" P
 Set A

$P \in A$ Yes

$P \notin A$ No

On Quant. Mec.



State $\xrightarrow{\text{Experiment}}$ $[\Psi]$ subspace

Take $\Psi_0 \in [\Psi]$; $|\Psi| = 1$

$$\Psi_0 = a\Psi_{\parallel} + b\Psi_{\perp} \quad (|\Psi_{\parallel}| = |\Psi_{\perp}| = 1)$$

$$|a|^2 + |b|^2 = 1$$

"Probability" that answer is Yes
 is equal to $P = |a|^2$

If we perform N experiments
with the same data then

in pN experiments answer "Yes"

in $(1-p)N$ "—————" "No"

Classical ————— Quantum Mech.

- or Logic.

Set

Question

subspace $V \subset H$

State.

a point

result of experim.

ray in H .

\in set "yes"

$\psi \in V$ "yes"

point

\notin set 'no'

$\psi \perp V$ "no"

$\Delta \psi, V \neq 0, \frac{1}{2}$ yes or no.

Set of questions is Set of questions

Distributive lattice

Modular lattice

$$A \cap (B \cup C) = A \cap B \cup A \cap C$$

$$x \leq b \Rightarrow$$

$(A \cap (B \cup C) \supset A \cap B \cup A \cap C)$
for arbitrary

$$x \vee (A \wedge B) \geq (x \vee A) \wedge B$$

Classical Mech.

Newton equations

$$\vec{F} = m \vec{a}$$

$$F = - \frac{\partial U}{\partial x^i}$$

$$m \ddot{x}^i = - \frac{\partial U}{\partial x^i}$$

Invariant of
 $SO(3)$

Quantum Mech.

Lagrangian

$$L = \frac{m \dot{x}^2}{2} - U(x)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}^i} = \frac{\partial L}{\partial x^i} \quad (\text{Euler-Lagrange})$$

Invariant of all
diffeomorphisms.

$$\begin{aligned} x &= x(q, t) \\ \dot{x} &= x'(q) \dot{q} + \dot{x}_t \end{aligned}$$

$$S(x, T) = \inf_{x(t): x(0)=0, x(T)=\Gamma} \int_0^T L(x(t), \dot{x}(t), t) dt$$

 $S(x, T)$ - classical action,[for free particle $S = \frac{m x^2}{2t}$]Invariance of solutions. Eq. are defined by
variational principle.

$$S(x, t): \quad \frac{\partial S}{\partial t} + H(x, p) = 0, \quad p = \frac{\partial S}{\partial x}$$

 $H(x, p)$ - Hamiltonian.The role of L and H is revealed
in Quantum Mechanics.

Variational principle \Rightarrow Euler Lagrange eq.

$$x \rightarrow x + \varepsilon h$$

$$\begin{aligned} S(x + \varepsilon h) &= \int L(x + \varepsilon h, \dot{x} + \varepsilon \dot{h}) dt = \\ &= \underbrace{\int L(x, \dot{x}) dt}_{S[x_0]} + \varepsilon \int \left(\frac{\partial L}{\partial x} h + \frac{\partial L}{\partial \dot{x}} \dot{h} \right) dt = \\ &= S[x_0] + \underbrace{\varepsilon \frac{\partial L}{\partial \dot{x}} h \Big|_0^T}_0 + \varepsilon \int \left(\frac{\partial L}{\partial x} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} \right) h dt \end{aligned}$$

$$x_0: \quad \frac{\partial L}{\partial x_0} = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_0} \right)$$

Euler-Lagrange equations do not depend on coordinates.

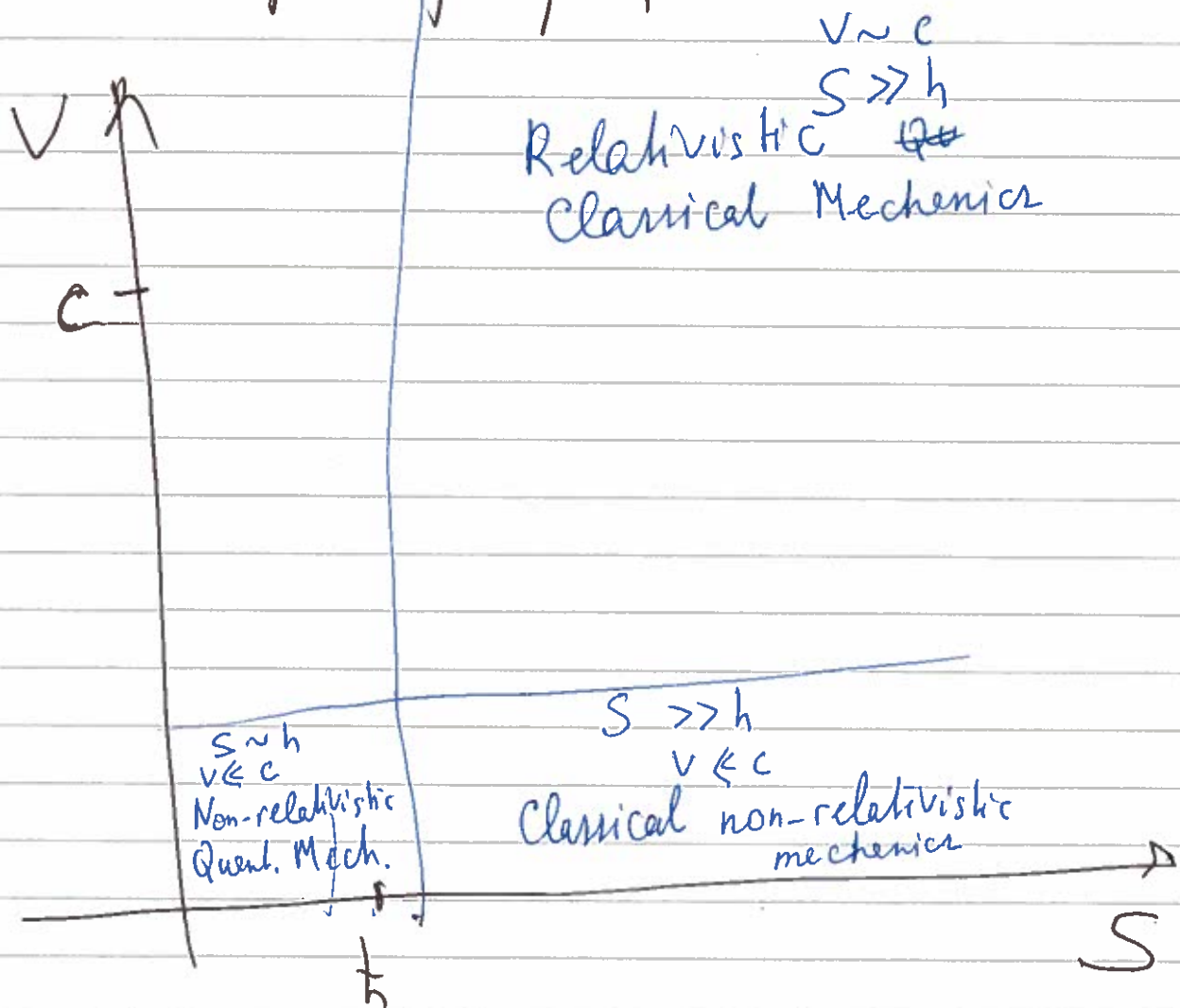
$$\frac{iS}{\hbar}$$

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$$\Psi \sim e$$

$$\Psi = \int e^{\int L(x, \dot{x}) dt} \mathcal{D}x(t)$$

Map of physics



$$\hbar = 6.10^{-27} \text{ J}\cdot\text{s}$$

$$1 \text{ J} = 10^7 \text{ erg}$$