

15 November 2018

Sixth lecture

Today we consider Schrödinger equation for harmonic oscillator:

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H}\Psi, \hat{H} = \frac{\hat{P}^2}{2m} + \frac{m\omega^2 q^2}{2}$$

Why harmonic oscillator ???

(In quantum field theory it is the basic system.
Try to give hints, why it is so. . .)

Consider an arbitrary Hamiltonian
 $\hat{H} = \sum_{i=1}^N \frac{\hat{P}_i^2}{2m_i} + U(q_1, q_2, \dots, q_N)$

H describes N degrees of freedom (dyn. system)

$$\begin{cases} \frac{dp_i}{dt} = -\frac{\partial U}{\partial q^i} \\ \frac{dq^i}{dt} = p_i \end{cases}$$

$$\frac{d^2q^i}{dt^2} = -\frac{\partial U}{\partial q^i}$$

Newton law.

Consider potential $U = U(q_1, \dots, q_N)$

in a vicinity of stability point:

$$U = U_0 + \frac{1}{2} \sum U_{ik} (q^i - q_0^i) (q^k - q_0^k)$$

$$\left. \left(\frac{\partial U}{\partial q^i} \right) \right|_{q_0} = 0 \quad (\text{stationary point})$$

~~$q_i \rightarrow q_i - q_{i0}$~~

approxim akely:

$$U \approx U_0 + \frac{1}{2} A_{ik} (q^i - q_{i0}) (q^k - q_{k0})$$

$q^0 \rightarrow q^0 \rightarrow q_{i0}^i$

We come to

$$H = \sum_{i=1}^N \frac{p_i^2}{2m_i} + \frac{1}{2} \sum A_{ijk} q^i q^k$$

Consider arbitrary linear transform, $q^i \rightarrow L_k^i q^k$
such that It does not change $(\dot{q}_1^2 + \dots + \dot{q}_n^2)$

$$\downarrow \\ L^\dagger L = 1 \Rightarrow L \text{ is orthogonal matrix}$$

one can find new coordinates such that

$$H = \sum \frac{p_i^2}{2m_i} + \frac{1}{2} \sum \lambda_i (q^i)^2 \quad (\text{Linear algebra !!})$$

all $\lambda_i \geq 0$ (if $\lambda_i < 0$ this is

non-stable equilibrium)

denote $A_i = \frac{m_i \omega_i^2}{2}$

$$H = \sum \frac{p_i^2}{2m_i} + \frac{1}{2} \sum \frac{m_i \omega_i^2 q_i^2}{2} = \quad (*)$$

$$= \sum_{i=1}^N \left(\frac{p_i^2}{2m_i} + \underbrace{\frac{m_i \omega_i^2 q_i^2}{2}}_{\text{free oscillator}} \right)$$

We see that (*) describes N
free (non-interacting) oscillators

Now look in detail oscillator:

$$\hat{H} = \frac{\hat{P}^2}{2} + \frac{m\omega^2 q^2}{2}$$

In coord. present: $\Psi = \Psi(q)$, $\hat{P} = i\hbar \frac{\partial}{\partial q}$, $\hat{q} = q$
 $i\hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi = \left(-\frac{\hbar^2}{2m} \Psi'' + \frac{m\omega^2 q^2}{2} \Psi \right)$

Find spectrum of this operator

$$\hat{H} \Psi_n = E_n \Psi_n \quad \left(\Psi(q, t) = \sum_n e^{\frac{i}{\hbar} E_n t} \Psi_n(q) \right)$$

Choose new coordinates ("kill dimension")

$$q = \sqrt{\frac{\hbar}{m\omega}} x, \quad \frac{d}{dq} = \frac{dx}{dq} \frac{d}{dx} = \sqrt{\frac{m\omega}{\hbar}} \frac{d}{dx}$$

$$\left(-\frac{\hbar^2}{2m} \Psi_n'' + \frac{m\omega^2 q^2}{2} \Psi_n \right) = E_n \Psi_n$$

$$\boxed{\cancel{\frac{\hbar^2}{2m} \left(x^2 - \frac{d^2}{dx^2} \right)} \Psi_n(x) = \frac{E}{\hbar\omega} \Psi_n}$$

We look carefully on this equation

$$\hat{a} = \frac{1}{\sqrt{2}} \left(x + \frac{d}{dx} \right)$$

$$\hat{a}^+ = \frac{1}{\sqrt{2}} \left(x - \frac{d}{dx} \right)$$

$$[\hat{a}, \hat{a}^+] = \frac{1}{2} \left[\left(x + \frac{d}{dx} \right), \left(x - \frac{d}{dx} \right) \right] =$$

$$\frac{1}{2} \left[\frac{d}{dx}, x \right] - \frac{1}{2} \left[x, \frac{d}{dx} \right] = 1.$$

$$\hat{a}^\dagger \hat{a} + \frac{1}{2} = \frac{1}{2} \left(x - \frac{d}{dx} \right) \left(x + \frac{d}{dx} \right) + \frac{1}{2} =$$

$$= \frac{1}{2} \left(x^2 - \frac{d^2}{dx^2} \right). \quad (\star)$$

Our equation

$$\frac{1}{2} \left(x^2 - \frac{d^2}{dx^2} \right) \psi_n = \frac{E_n}{\hbar \omega} \psi_n \iff$$

$$\begin{cases} (\hat{a}^\dagger \hat{a} + \frac{1}{2}) |\psi\rangle = \epsilon |\psi\rangle \\ [\hat{a}, \hat{a}^\dagger] = 1 \end{cases}$$

$$\text{In fact } \hat{a} = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{q} + \frac{i}{m\omega} \hat{p} \right)$$

$$\hat{a}^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{q} - \frac{i}{m\omega} \hat{p} \right)$$

These operators are called
creation (\hat{a}^+)
annihilation (\hat{a})

Operators

$$\hat{H} = \hat{Q}^\dagger Q = \frac{m\omega}{2\hbar} \left(\hat{q} - \frac{i}{m\omega} \hat{p} \right) \left(\hat{q} + \frac{i}{m\omega} \hat{p} \right) =$$

$$= \left(\frac{m\omega^2 q^2}{2} + \frac{\hat{p}^2}{2m} \right) \hbar\omega - \frac{1}{2}$$

↓

$$\hat{H} = \hbar\omega \left(\hat{Q}^\dagger \hat{Q} + \frac{1}{2} \right), \quad [\hat{Q}, \hat{Q}^\dagger] = L$$

This ~~is~~ representation is very convenient.

Find solutions:

$$\hbar\omega (\hat{Q}^\dagger \hat{Q} + \frac{1}{2}) |\psi_n\rangle = E_n |\psi_n\rangle$$

Consider ϕ : $\hat{Q}^\dagger \phi = 0$

ϕ is called
Vacuum

$$\begin{aligned} \hat{Q}^\dagger \phi &= 0 \\ \frac{1}{\sqrt{2}} \left(\frac{d}{dx} + x \right) \phi &= 0 \\ \phi &= C e^{-\frac{x^2}{2}} \end{aligned}$$

$$\phi = |0\rangle$$

Denote $\phi = |0\rangle$

(if ϕ is just notation)

$$\hat{H}|0\rangle = \left(\hat{a}^\dagger \hat{a} + \frac{1}{2}\right)|0\rangle = \frac{1}{2}|0\rangle$$

since $\hat{a}|0\rangle = 0$

Annihilation operator "kills" vacuum.

Consider

$$|1\rangle = \hat{a}^\dagger |0\rangle \text{ . Use } [\hat{a}, \hat{a}^\dagger] = 1 \Rightarrow$$

$$\begin{aligned} \hat{H}|1\rangle &= \left(\hat{a}^\dagger \hat{a} + \frac{1}{2}\right) \hat{a}^\dagger |0\rangle = \\ &= \left\{ \hat{a}^\dagger \left(\hat{a}^\dagger \hat{a} + 1 \right) + \frac{1}{2} \hat{a}^\dagger \right\} |0\rangle = \\ &= \left[\hat{a}^\dagger \hat{a} + \frac{3}{2} \hat{a}^\dagger \right] |0\rangle = \end{aligned}$$

$$= \frac{3}{2} \underbrace{|\hat{a}^\dagger |0\rangle}_{|1\rangle} = \frac{3}{2} |1\rangle$$

Energy of the state $|1\rangle$ is equal to $\boxed{\frac{3}{2}}$

Show! — Is it true that $(|1\rangle \neq 0)$

Check it

$$\langle 1 | 1 \rangle = \langle \hat{a}^\dagger | 0 \rangle, \quad \hat{a}^\dagger | 0 \rangle =$$

$$= \langle 0 | \hat{a} \hat{a}^\dagger | 0 \rangle = \langle 0 | \hat{a}^\dagger \hat{a} + 1 | 0 \rangle =$$

$$= 1 \neq 0$$

In these calculations we just used
 $\hat{a}|0\rangle = 0$.

annihilation operator kills the vacuum

$$\hat{a}\hat{a}^\dagger - \hat{a}^\dagger\hat{a} = 1 \quad ([\hat{a}, \hat{a}^\dagger] = 1)$$

Commutation relations

Lemma. $[\hat{H}, \hat{a}] = -\hat{a}$

$$[\hat{H}, \hat{a}^\dagger] = \hat{a}^\dagger$$

Check it!

Use it.

$$\text{Let } \Psi: \hat{H}\Psi = E\Psi$$

$$\hat{H}(\hat{a}^\dagger\Psi) = (\hat{a}^\dagger\hat{H} + \hat{a}^\dagger)\Psi = (E+1)\hat{a}^\dagger\Psi$$

We proved:

$\hat{H}\Psi$ has energy $E \implies \hat{a}^\dagger\Psi$ has energy $E+1$

!!!

(It has to be checked that $\hat{a}^\dagger\Psi \neq 0$:

$$(\langle \hat{a}^\dagger\Psi, \hat{a}^\dagger\Psi \rangle = \langle \Psi, \hat{a}\hat{a}^\dagger\Psi \rangle = 1 - |\hat{a}\Psi|^2 \geq 1)$$

Check it!

(See exercises in the homework 6)

- 8

$$|n\rangle = \frac{(a^+)^n}{\sqrt{n!}} |0\rangle$$

$$\langle n|m\rangle = \begin{cases} 1 & \text{if } n=m \\ 0 & \text{if } n \neq m. \end{cases}$$

$$\hat{a}^+|n\rangle = \sqrt{n+1} |n+1\rangle$$

$$\hat{a}|n\rangle = \sqrt{n} |n-1\rangle$$

\hat{a}, \hat{a}^+ -

\hat{a}^+ transforms

k -th level of

oscillator to $(k+1)$ -th
level

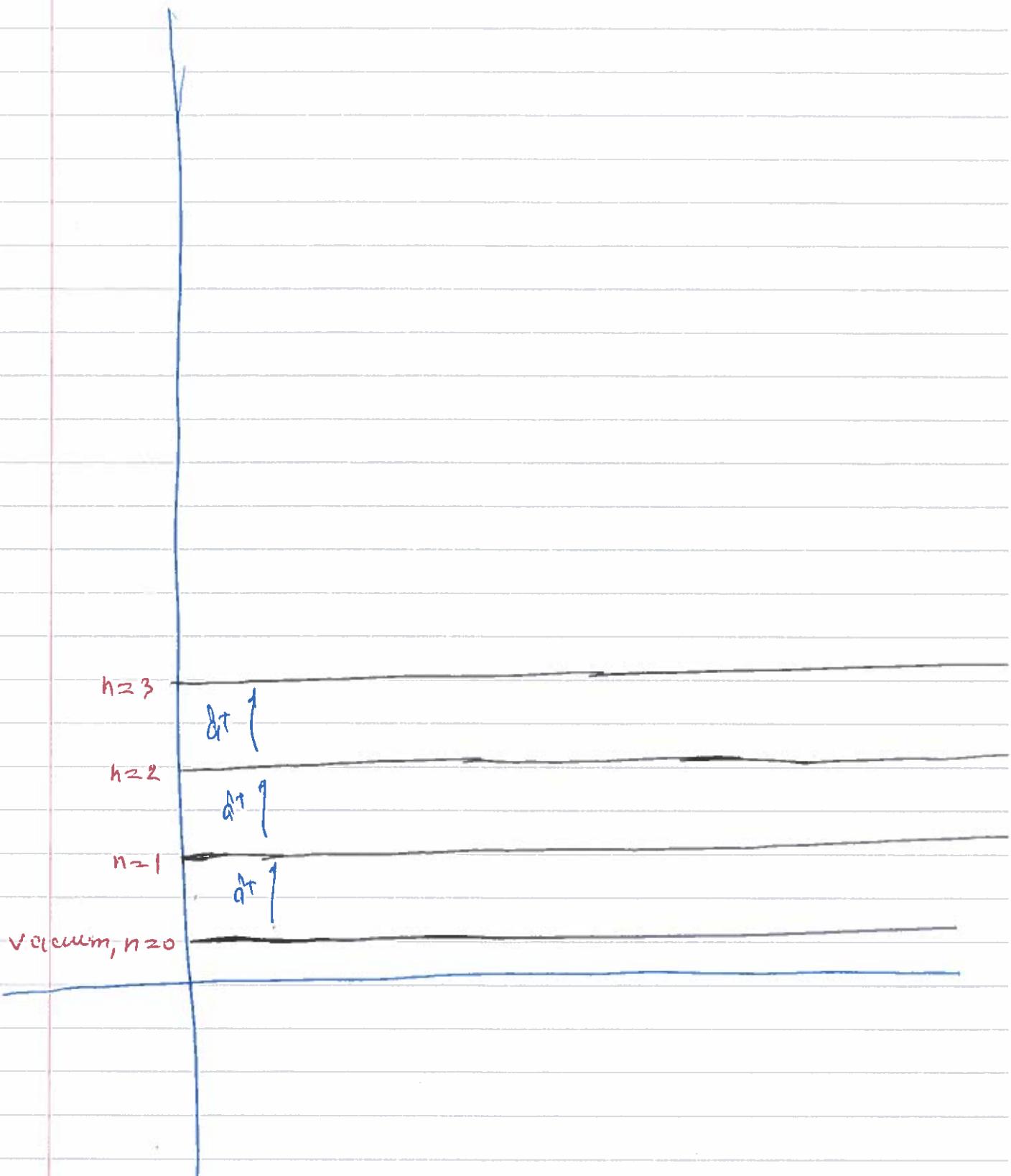
\hat{a} : k -th level $\Rightarrow (k-1)$ th
level,

\hat{a}^+ add per tide
with Energy $\hbar\omega$

\hat{a} 'kills' particle
with energy $\hbar\omega$

We see that

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Fock LADDER